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Upconversion phenomena in the Yb³⁺ doped transition metal compounds Rb₂MnCl₄ and CsMnBr₃

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Abstract

Crystals of Yb³⁺ doped Rb₂MnCl₄ and CsMnBr₃ show orange-red Mn^{2+ 4}T_{1g} \rightarrow ⁶A_{1g} upconverted luminescence under near-infrared Yb³⁺ excitation. By short 10 ns pulse excitations the upconversion (UC) process for both compounds has been determined to consist of a sequence of ground and excited-state absorption steps. This mechanism can gain nonzero probability by assuming an exchange mechanism between Yb³⁺ and Mn²⁺. The efficiency of the UC process is dependent on the structural nature of the Yb³⁺–Mn²⁺ connection. It is three orders of magnitude more efficient in Rb₂MnCl₄: Yb³⁺ with a corner-sharing Yb³⁺–Cl–Mn²⁺ arrangement than in the face-sharing Yb³⁺–Br₃–Mn²⁺ arrangement of CsMnBr₃: Yb³⁺. © 2001 Elsevier Science B.V. All rights reserved.

Keywords: Upconversion; Exchange interaction

1. Introduction

Upconversion (UC) is established in rare earth compounds as well as in a small number of transition metal ion systems [1,2]. A new type of UC process has recently been observed in mixed Yb^{3+}/Mn^{2+} systems which includes cooperative electronic transitions based on exchange interactions between the two ions [3–5]. In the following, we describe the two title compounds which differ in their bridging geometry, ligand field strength and magnetic behavior. We are interested in the underlying upconversion mechanisms and the trends in UC efficiency when the chemical environment and the bridging geometry is changed.

2. Experimental section

Single crystals of Rb_2MnCl_4 : Yb^{3+} and $CsMnBr_3$: Yb^{3+} were grown by the Bridgman technique using stoichiometric amounts of the corresponding chlorides and bromides, respectively, and a trace of Yb^{3+} , as described in Ref. [6]. The real ytterbium concentration is 0.5 ± 0.1 mol% measured by ICP-OES for the CsMnBr₃: Yb^{3+} compound. From absorption measurements it was concluded, that the concentration of Yb^{3+} in Rb_2MnCl_4 : Yb^{3+} is <0.1%. The spectroscopic measurements were carried out as described in Ref. [4].

3. Results

In both crystals near-IR (NIR) excitation into a ${}^2F_{7/2} \rightarrow {}^2F_{5/2}$ absorption band of Yb³⁺ leads to

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orange-red broadband UC luminescence at 10 K, which is readily assigned to the $Mn^{2+} {}^{4}T_{1g} \rightarrow {}^{6}A_{1g}$ transition. Fig. 1 shows the corresponding optical spectroscopic results for $Rb_2MnCl_4: Yb^{3+}$. In Fig. 1b both the downconverted $Yb^{3+} {}^{2}F_{5/2} \rightarrow {}^{2}F_{7/2}$ luminescence consisting of sharp lines around 10,000 cm⁻¹ and the upconverted broadband luminescence centered around 16,000 cm⁻¹ are seen. Note the different scales of the two spectra. For a laser power of 140 mW focused with a 53 mm focal lens on the crystal the ratio of VIS : NIR emitted photons is 30% at 15 K. Fig. 1a shows the excitation spectrum in the NIR of the UC luminescence. It correspond to the ${}^{2}F_{7/2} \rightarrow {}^{2}F_{5/2}$ absorption bands.

Fig. 2a shows the σ -polarized absorption spectrum of CsMnBr₃: Yb³⁺ in the NIR and VIS. Both the ${}^{2}F_{7/2} \rightarrow {}^{2}F_{5/2}$ absorption lines above 10,000 cm⁻¹ and the first d–d bands of CsMnBr₃: Yb³⁺ are recognized. Fig. 2b shows the downconverted NIR luminescence of Yb³⁺ and the upconverted luminescence at 12 K after Yb³⁺ excitation around 10,600 cm⁻¹. In this figure the scaling factor is three orders of magnitude

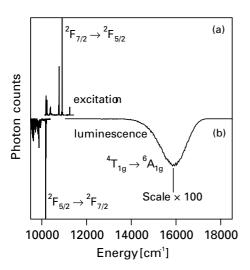


Fig. 1. (a) 15 K excitation spectrum monitoring the orange $Mn^{2+4}T_{1g} \rightarrow {}^{6}A_{1g}$ upconversion luminescence centered at 16,000 cm⁻¹. (b) Survey luminescence and upconversion luminescence spectra at 15 K of Rb₂MnCl₄: Yb³⁺ (<0.1%) excited with E//c at 10,777 cm⁻¹. Note the magnification of the upconversion by a factor of 100.

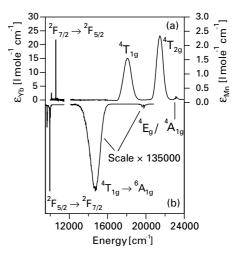


Fig. 2. (a) $E \perp c$ polarized absorption spectrum of CsMnBr₃: Yb³⁺ (0.5%) at 12 K. (b) 12 K luminescence and upconversion spectra after NIR excitation at 10,596 cm⁻¹ showing the Yb³⁺ luminescence around 10,000 cm⁻¹, the red Mn²⁺⁴T_{1g} \rightarrow ⁶A_{1g} upconverted luminescence centered at 14,500 cm⁻¹ and the green cooperative Yb³⁺ pair UC luminescence around 20,000 cm⁻¹. The UC luminescences are scaled up by a factor of 135,000.

larger than in Fig. 1, indicating that for equivalent excitation power at 12 K the ratio of VIS : NIR photons is three orders of magnitude smaller in this crystal than in Rb_2MnCl_4 : Yb^{3+} . In Fig. 2b we also recognize a very weak structured luminescence band around 20,000 cm⁻¹. This is assigned to a cooperative UC luminescence of Yb^{3+} pairs, which are present in small concentration in this crystal. This type of UC process has been described in numerous Yb^{3+} compounds [7]. We note that in CsMnBr₃ : Yb^{3+} this luminescence can only be observed because it happens to lie in the absorption window between the ${}^{4}T_{1g}$ and ${}^{4}T_{2g}$ absorption bands of Mn²⁺.

Fig. 3 shows the time dependence at 12 K of the upconverted luminescence intensity in a semilogarithmic representation after a 10 ns laser pulse into the Yb³⁺ absorption in the NIR for both crystals. We notice that the curves have no rise and correspond to single exponential decays. Decay times of 9.8 ms and 337 μ s are derived for Rb₂MnCl₄ : Yb³⁺ and CsMnBr₃ : Yb³⁺ at 12 K, respectively.

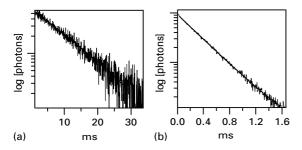


Fig. 3. Temporal behavior of the upconverted Mn^{2+} $^4T_{1g} \rightarrow ^6A_{1g}$ 12 K luminescence intensity exciting with a 10 ns laser pulse at 10,201.7 cm $^{-1}$ and at 10,596 cm $^{-1}$ in (a) Rb₂MnCl₄: Yb $^{3+}$ and (b) CsMnBr₃: Yb $^{3+}$, respectively. The straight lines show the single exponential fits to the experimental data.

4. Discussion

Fig. 4 shows the relevant structure elements of the two lattices. Rb_2MnCl_4 crystallizes in a layertype perovskite structure with corner-sharing $[MnCl_6]^{4-}$ octahedra in two dimensions perpendicular to the tetragonal *c*-axis, see Fig. 4a. It is not known how Yb^{3+} is incorporated and charge compensation is achieved in this lattice. From absorption experiments we know that our crystal contains less than 0.1% Yb^{3+} , even though 1% of Yb^{3+} was added to the melt. Incorporation of Yb^{3+} and charge compensation is evidently difficult in this system. The situation is

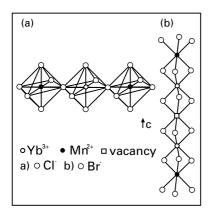


Fig. 4. Relevant elements of the Rb_2MnCl_4 : Yb^{3+} and $CsMnBr_3$: Yb^{3+} structures representing schematically the corner-sharing (a) and face-sharing (b) connectivity between the Mn^{2+} ions.

very different in CsMnBr₃, which crystallizes in a one-dimensional structure with face-sharing $[MnBr_6]^{4-}$ octahedra forming chains along the hexagonal *c*-axis. In this type of lattice the incorporation of trivalent ions such as Yb³⁺ is much easier and well established [8]. As shown in Fig. 4b, the most likely arrangement is one in which a sequence Yb³⁺-vacancy-Yb³⁺ is replacing three Mn²⁺ ions. In this crystal the actual Yb³⁺ concentration of 0.5% is close to the one used in the melt for the crystal growth.

In the following, we will briefly discuss what we believe to be the main mechanism for the unusual UC phenomenon. In addition the question of the very different UC efficiencies in the two title systems will be addressed. Similar phenomena have very recently been reported for $RbMnCl_3$: Yb^{3+} and CsMnCl₃: Yb³⁺ [3]. In CsMnCl₃: Yb³⁺ the UC efficiency was found to be very high, similar to $Rb_2MnCl_4: Yb^{3+}$ reported here, whereas in $RbMnCl_3$: Yb³⁺ it is significantly lower. The studies on the title compounds were motivated by these differences. We suspect a correlation with the structure, in particular the bridging between the incorporated Yb^{3+} ions and its Mn^{2+} neighbors. That is why we chose lattices with different bridging geometries between adjacent Mn2+ ions in the present study: corner-sharing octahedra in Rb_2MnCl_4 : Yb³⁺ (Fig. 4a) and face-sharing octahedra in CsMnBr₃: Yb³⁺ (Fig. 4b).

From the UC excitation spectra in Figs. 1 and 2 it clearly follows that a ${}^{2}F_{7/2} \rightarrow {}^{2}F_{5/2}$ absorption is the first step in the upconversion process in both systems studied here. And the absence of a rise in the transient curves of Fig. 3 rules out an energy transfer UC mechanism. Both the ground-state absorption (GSA) step and the upconversion process occur within the 10 ns of the laser pulse. We conclude that the second step is an excitedstate absorption (ESA) step from the intermediate level, and the total process thus a GSA/ESA sequence. Whereas the GSA is a normal Yb^{3+} absorption, the ESA causes considerable conceptual problems. It corresponds to an electronic transition from an intermediate state localized on Yb^{3+} to a final state localized on Mn^{2+} . This requires some kind of coupling between Yb^{3+} and Mn^{2+} . This could be either an exchange or an

electric multipole coupling. We believe that it is based on exchange interactions, and the experimentally found strong dependence of the upconversion efficiency on the bridging geometry confirms this. We use a $Yb^{3+}-Mn^{2+}$ dimer picture to account for the observed effects. We are aware that this is approximate, because Rb₂MnCl₄: Yb³⁺ is magnetically ordered at 12 K [9] and, in addition, Yb^{3+} is likely to have more than one Mn^{2+} neighbor in Rb_2MnCl_4 : Yb^{3+} . But the dimer picture, which is shown in Fig. 5, does contain the essential features for an understanding. And since a great deal of work has been done on transition metal dimers, we can use some old concepts from these studies as a basis. As seen in Fig. 5 the dimer states have double labels, one part from each partner. And the second step in the upconversion process, which causes the complications in a picture with separate Yb^{3+} and Mn²⁺ chromophores, is quite a normal ESA step within the dimer. How does the ESA step gain intensity? We believe that it is the same mechanism which gives intensity to pair transitions of Cr^{3+} . which are forbidden in a single-ion mechanism [10]. Tanabe and coworkers proposed such a mechanism to account for the absorption spectra of magnetically coupled Mn^{2+} systems [11]. It has been shown that this mechanism can enhance the intensity of transitions which are spin-forbidden in the single ion, by several orders of magnitude. It can also account for the occurrence of so-called cooperative transitions, such as double excitations in a Cr³⁺ dimer, which cannot occur in single ions [12]. The

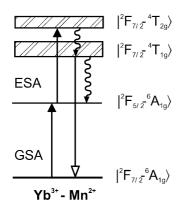


Fig. 5. Schematic representation of the UC mechanism in a Yb^{3+} – Mn^{2+} dimer picture. Full arrow lines and curly arrows represent radiative and nonradiative relaxation processes.

ESA step in Fig. 5 is such a cooperative electronic transition, which takes the dimer from a nominally Yb^{3+} centered excited state to a nominally Mn^{2+} excited state.

For an exchange mechanism we expect a dependence of the efficiency of the process on the chemical and structural nature of the Yb³⁺-Mn²⁺ connection. For an electric multipole mechanism, on the other hand, we would expect a strong decrease of the efficiency with increasing Yb^{3+} - Mn^{2+} distance. The present study very clearly shows a correlation with the nature of the Yb^{3+} - Mn^{2+} bridge. If the $Yb^{3+}-Mn^{2+}$ distance was important we would expect a higher efficiency of the UC process in CsMnBr₃: Yb^{3+} , because in this lattice the expected Yb³⁺–Mn²⁺ distance is 3.26 Å, much shorter than in Rb_2MnCl_4 : Yb^{3+} with 5.05 Å. It appears safe to conclude that the corner-sharing Yb^{3+} -Cl-Mn²⁺ arrangement is more efficient than the face-sharing Yb³⁺-Br₃- Mn^{2+} in promoting the cooperative ESA step in Fig. 5. This conclusion is in very good agreement with earlier suspicions based on RbMnCl₃: Yb^{3+} and $CsMnCl_3$: Yb³⁺. But it is only by the type of study presented here, i.e. a systematic variation of the bridging geometries, that we arrive at unambiguous conclusions.

In all the Yb³⁺ doped Mn²⁺ lattices studied so far this unusual UC is quenched at room temperature. We ascribe this to the occurrence of a nonradiative Mn²⁺ to Yb³⁺ multiphonon relaxation process which becomes efficient at high temperatures. Evidence for this is the observation of increasing Yb³⁺ luminescence with increasing temperature upon Mn²⁺ excitation in the VIS. Unfortunately, it is hard to circumvent this intrinsic loss channel by clever chemical engineering.

Acknowledgements

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